

## Limit theorems for age-dependent branching process with emigration

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### Abstract

In this paper we consider age-dependent branching processes with emigration as a part of the same process. We prove limit theorems for the process that describe number of particles which emigrated.

**Keywords.** branching processes, stochastic additive functional, critical branching processes, emigration, perron root, moments.

### 1. Introduction

We will define multi-type age –dependent branching process with emigration by analogy with [2].

As in the multi-type process without emigration we have  $n$  types of particles  $T_1, T_2, \dots, T_n$ . Each  $T_i$ -th type particle has random duration of ‘presence’ in process  $\tau_i$  with distribution function

$$P(\tau_i \leq t) = G^i(t), G^i(0+) = 0$$

We will assume that  $G^i(t)$  are absolutely continuous.

At the end of its presence particle of any type is transformed into arbitrary number of particles of any types or emigrate. And let's define type  $T_0$  as a type of particles which emigrated. Conditional probability (if transformation took place when the age attained by the original particle was  $u$ )  $p_\alpha^i(u)$  of transformation into a set consisting of  $\alpha_i$   $T_i$  – th type particles,  $i = \overline{0, n}$ , where  $\alpha = (\alpha_0, \alpha_1, \alpha_2, \dots, \alpha_n)$  is  $n+1$ -dimensional vector and  $\alpha_0$  can be 0 ( in this case other components of vector take non-negative integer values) or 1 ( in this case other components of vector are zeroes). Evolution of particle is defined by joint distribution of random variable  $\tau_i$  and random vector  $v_i = (v_i^0, v_i^1, \dots, v_i^n)$ , which characterize progeny of this particle.

$$P(\tau_i \in B, v_i = \alpha) = \int_B p_\alpha^i(u) dG^i(u)$$

Particle of type  $T_0$  at any moment transforms into itself (or doesn't transform, you can interpret it either way), so it can be viewed as  $T_0$  -th type particle has degenerate at zero life duration distribution function ( $G^0(t) = 1, \forall t > 0$ ). Also we will assume that  $p_1^0(u) = 1, p_\alpha^0(u) = 0$  for all  $u \geq 0$  and all vectors  $\alpha \neq (1, 0, \dots, 0)$ , where  $1 = \underbrace{(1, 0, \dots, 0)}_{n+1}$ . Vector  $\mu_i(t) = (\mu_i^0(t), \mu_i^1(t), \dots, \mu_i^n(t))$  denotes number of particles of

types  $T_0, T_1, T_2, \dots, T_n$  at the moment  $t$ , under the condition, that at initial moment there existed one  $T_i$  – th type particle. And let's also assume, that vector are right continuous.

Let's denote by  $P^i(*)$  conditional probability, under condition that at initial moment of time there existed one particle of type  $T_i$ .

Let's also introduce generating functions

$$h^i(t, s) = \sum_\alpha p_\alpha^i(t) s^\alpha \text{ and } F^i(t, s) = \sum_\alpha P^i(\mu_i(t) = \alpha) s^\alpha, i = \overline{0, n},$$

$$s = (s_0, s_1, \dots, s_n), s^\alpha = s_0^{\alpha_0}, s_1^{\alpha_1}, \dots, s_n^{\alpha_n},$$

$$F(t, s) = (F^0(t, s), F^1(t, s), \dots, F^n(t, s)), h(t, s) = (h^0(t, s), h^1(t, s), \dots, h^n(t, s)).$$

It's clear that  $F^0(t, s) = h^0(t, s) = s_0$  for all  $t$ .

**Lemma 1.** Generating functions  $F^i(t, s)$  satisfy with  $s \leq 1$  and  $t \geq 0$  next system of integral equations:

$$F^i(t, s) = \int_0^t h^i(u, F(t-u, s)) dG^i(u) + s_i(1 - G^i(t)), \quad i = \overline{0, n}. \quad (1)$$

Proof of this lemma is analogical to proof of the theorem 1 [2], p.234.

## 2. Moments

During the paper we will use following notations:

$$\begin{aligned} \alpha_j^i(u) &= \left. \frac{\partial h^i(u, s)}{\partial s_j} \right|_{s=1}, \quad A_j^i = \int_0^\infty \alpha_j^i(u) dG^i(u), \\ b_{jk}^i(u) &= \left. \frac{\partial h^i(u, s)}{\partial s_j \partial s_k} \right|_{s=1}, \quad B_{jk}^i = \int_0^\infty b_{jk}^i(u) dG^i(u), \\ A_{\rho k}^l &= \int_0^\infty u e^{-\rho u} a_k^l(u) dG^l(u), \quad B_{\rho jk}^l = \int_0^\infty e^{-2\rho u} b_{jk}^l(u) dG^l(u), \\ M^k &= \int_0^\infty u dG^k(u), \quad M_a^l = \int_0^\infty u a_k^l(u) dG^l(u), \\ B &= \sum_{l, k, m=1}^n B_{mk}^l v^l u^k u^m, \quad M_a = \sum_{l, k=1}^n M_a^l v^l u^k, \quad M_0 = \sum_{l=1}^n v^l \int_0^\infty a_0^l(u) dG^l(u), \end{aligned}$$

where  $\rho$  denotes such number, that perron root of matrix  $e^{-\rho t} \|A_j^i\|_{i, j=1, \overline{n}}$  equals to one,  $u_\rho = (u_\rho^1, u_\rho^2, \dots, u_\rho^n)$  and  $v_\rho = (v_\rho^1, v_\rho^2, \dots, v_\rho^n)$  denote right and left eigenvectors of this matrix,  $(v_\rho, 1) = \sum_{k=1}^n v_\rho^k = 1$  (in critical case  $\rho = 0$  and we denote  $u_\rho = u, v_\rho = v$ ).

We will call described above branching process (b.p)  $\xi$ . We will also introduce additional b.p.  $\xi'$ , which differ from  $\xi$  only by the fact, that conditioned probability of transforming into the empty set  $p_0^i(u)$  equals to the sum of  $p_0^i(u)$  and  $p_1^i(u)$  where  $0 = \underbrace{(0, 0, \dots, 0)}_{n+1}$ . So process  $\xi'$  could be seen simply as a multi-type age dependent process. We will call process  $\xi$  subcritical (critical or supercritical) iff  $\xi'$  is subcritical (critical or supercritical) (see, for example, [2]). It's clear that moments  $A_j^{i'}$  of process  $\xi'$ , for  $i, j = \overline{1, n}$ , and it's also clear that  $a_0^i(u) = p_1^i(u)$ .

Now consider matrix  $A = \|A_j^i\|_{i, j=0, \overline{n}}$  of moments of process  $\xi$

Obviously  $A_0^0 = 1$  and  $A_j^0 = 0 \quad \forall j = \overline{1, n}$ . Then has  $A$  from  $\begin{bmatrix} 1 & 0 \\ A_{10} & A' \end{bmatrix}$ , where  $A'$  is matrix of moments of process  $\xi'$ ,  $A_{10} = (A_0^1, A_0^2, \dots, A_0^n), 1 = 1, 0 = \underbrace{(0, 0, \dots, 0)}_n$ .

## 3. Supercritical case.

As we mentioned above, our b.p. is a special type of decomposable branching processes, but it turns out that we could apply theory of stochastic additive functionals from b.p. (see [3]) to our case. In fact we

can interpret processes  $\mu_0^i(t)$  where amount of product  $\gamma(t)$  (which depends on type of particle, age at the moment  $t$  and life duration (we called it presence above)), produced by 1 particle during the time period  $0 \leq t \leq \tau$ , ( $\tau$  as usually denote life duration) is

$$\gamma(t) = \begin{cases} 1, & \text{if particle emigrates at moment } \tau, \\ 0, & \text{otherwise,} \end{cases}$$

Analogy of theorem 2 (and corollary 1) [3] takes place, which also can be proven using slight modification of results obtained in [1], if we approach this as decomposable b.p.:

**Theorem 1.** If process  $\xi$  is supercritical,  $A_{\rho k}^l(u)$  and  $B_{\rho k}^l$  are finite,  $i, j, k = \overline{1, n}$ , than random variables (r.v.)  $\frac{e^{-\rho t} \mu_0^i(t)}{K}, \frac{e^{-\rho t} \mu_1^i(t)}{K_1}, \dots, \frac{e^{-\rho t} \mu_n^i(t)}{K_n}$  converge as  $t \rightarrow \infty$  in square mean to the same limit  $\mu^i$ , where

$$K = \frac{\sum_{l=1}^n \nu_\rho^l \int_0^\infty e^{-\rho t} \int_0^t a_0^i(u) dG^i(u) dt}{\sum_{l,m=1}^n \nu_\rho^l u_\rho^m \int_0^\infty t e^{-\rho t} a_j^i(t) dG^l(t)}, \quad K_l = \frac{\nu_\rho^l \int_0^\infty e^{-\rho t} \int_0^t (1 - G^i(t)) dt}{\sum_{l,m=1}^n \nu_\rho^l u_\rho^m \int_0^\infty t e^{-\rho t} a_j^i(t) dG^l(t)}, \quad l = \overline{1, n}$$

Furthermore, Laplace transform  $\phi^l(s)$  of limit r.v.  $\mu^l$ , satisfies equation

$$\phi^l(s) = \int_0^\infty h^l(u, \phi(se^{-\rho u})) dG^l(u), \quad (2)$$

with initial condition  $\phi^l(0) = u_\rho$ , where  $\phi(s) = (\phi^1(s), \dots, \phi^n(s))$ .

**Corollary 1.** If conditions of theorem 1 are satisfied, than distributions

$$P^i \left( \frac{\mu_0^i(t)}{\sum_{k=1}^n \mu_k^i(t)} \leq x \mid \sum_{k=1}^n \mu_k^i(t) > 0 \right), \quad i = \overline{1, n}$$

converge to the degenerate distribution, localized at the point

$$\frac{\sum_{l=1}^n \nu_\rho^l \int_0^\infty e^{-\rho t} \int_0^t a_0^i(u) dG^i(u) dt}{\sum_{l=1}^n \nu_\rho^l \int_0^\infty e^{-\rho t} (1 - G^i(t)) dt}.$$

#### 4.Critical case.

In order to formulate theorem for critical case we will compare processes  $\mu_0^i(t)$  with processes  $N^i(t)$  – total number of particles born by the moment of time  $t$ , if at the moment  $t=0$  there existed one particle of type  $T_i$ .

Let  $N_j^i(t)$  denote number of particles of type  $T_j$ , born by  $t$ , the  $N^i(t) = \sum_{j=1}^n N_j^i(t)$ . It is known [4] (in this paper is considered case with constant  $p_\alpha^i(u) = p_\alpha^i$  for all  $u$ , but

$$E\left(\exp\left\{i\sum_{j=1}^n \beta^j N_j^i(t)/v^j t^2\right\} \mid \sum_{j=1}^n \mu_j^i(t)\right) \xrightarrow{t \rightarrow +\infty} \\ \left(2B\sum_{j=1}^n \beta^j\right)^{1/2} / M_\alpha \left( \operatorname{sh}\left(\left(2B\sum_{j=1}^n \beta^j\right)^{1/2} / M_\alpha\right) \right).$$

Then by letting  $\beta^j = v^j \beta$ , and since  $\sum_{j=1}^n v^j = 1$ , we get

$$E\left(\exp\left\{i\beta^j N^i(t)/t^2\right\} \mid \sum_{j=1}^n \mu_j^i(t)\right) \xrightarrow{t \rightarrow +\infty}, \quad \Leftrightarrow (2B\beta)^{1/2} / M_\alpha \left( \operatorname{sh}\left((2B\beta)^{1/2} / M_\alpha\right) \right). \quad (3)$$

Also from [4] we can establish asymptotic behavior of moments  $E^i(N_i(t))$  and  $E^i(N_i^2(t))$ :

$$E^i(N_i(t)) \sim \frac{u^i t}{M_\alpha}, \quad E^i(N_i^2(t)) \sim B \frac{u^i t^3}{3(M_\alpha)^3}.$$

**Theorem 2.** If the following conditions are satisfied:

- i) integrals  $M^j, M_a^{jk}, B_{jk}^i$  are finite;
- ii)  $\int_0^\infty \int_t^\infty a_0^l(u) dG^l(u) dt = k_0^l < +\infty$ ;
- iii)  $\lim_{t \rightarrow \infty} t^2 \int_t^\infty a_l^k(u) dG^l(u) < +\infty, \quad \lim_{t \rightarrow \infty} t^2 (1 - G^l(t)) < +\infty, \quad l, k, j = \overline{1, n}$ , then

$$E\left(\exp\left\{i\beta^j \mu_0^i(t)/t^2\right\} \mid \sum_{j=1}^n \mu_j^i(t)\right) \xrightarrow{t \rightarrow +\infty}, \quad (2B M_0 \beta)^{1/2} / M_\alpha \left( \operatorname{sh}\left((2B M_0 \beta)^{1/2} / M_\alpha\right) \right)$$

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